Magnetic Boundary Conditions

Problem/3< The *x*-*y* plane separates two magnetic media with magnetic permeabilities μ_1 and μ_2 , as shown in Fig.3. If there is no surface current at the interface and the magnetic field in medium 1 is

$$\mathbf{H}_1 = \hat{\mathbf{x}} H_{1x} + \hat{\mathbf{y}} H_{1y} + \hat{\mathbf{z}} H_{1z},$$

find:

- (a) **H**₂,
- (**b**) θ_1 and θ_2 , and
- (c) evaluate \mathbf{H}_2 , θ_1 , and θ_2 for $H_{1x} = 2$ (A/m), $H_{1y} = 0$, $H_{1z} = 4$ (A/m), $\mu_1 = \mu_0$, and $\mu_2 = 4\mu_0$.



Fig: Adjacent magnetic media.

Solution:

(a) From,

$$\mu_1 H_{1n} = \mu_2 H_{2n},$$

and in the absence of surface currents at the interface,

$$H_{1t} = H_{2t}$$
.

In this case, $H_{1z} = H_{1n}$, and H_{1x} and H_{1y} are tangential fields. Hence,

$$\mu_1 H_{1z} = \mu_2 H_{2z}, H_{1x} = H_{2x}, H_{1y} = H_{2y},$$

and

$$\mathbf{H}_2 = \hat{\mathbf{x}} H_{1x} + \hat{\mathbf{y}} H_{1y} + \hat{\mathbf{z}} \frac{\mu_1}{\mu_2} H_{1z}.$$

(b)

$$H_{1t} = \sqrt{H_{1x}^2 + H_{1y}^2},$$

$$\tan \theta_1 = \frac{H_{1t}}{H_{1z}} = \frac{\sqrt{H_{1x}^2 + H_{1y}^2}}{H_{1z}},$$

$$\tan \theta_2 = \frac{H_{2t}}{H_{2z}} = \frac{\sqrt{H_{1x}^2 + H_{1y}^2}}{\frac{\mu_1}{\mu_2}H_{1z}} = \frac{\mu_2}{\mu_1} \tan \theta_1.$$

(c)

$$\mathbf{H}_{2} = \hat{\mathbf{x}} 2 + \hat{\mathbf{z}} \frac{1}{4} \cdot 4 = \hat{\mathbf{x}} 2 + \hat{\mathbf{z}} \quad (A/m),$$

$$\theta_{1} = \tan^{-1} \left(\frac{2}{4}\right) = 26.56^{\circ},$$

$$\theta_{2} = \tan^{-1} \left(\frac{2}{1}\right) = 63.44^{\circ}.$$

Problem/4
Civen that a current sheet with surface current density $\mathbf{J}_s = \hat{\mathbf{x}} \, 8 \, (\text{A/m})$
exists at y = 0, the interface between two magnetic media, and $\mathbf{H}_1 = \hat{\mathbf{z}} \, 11 \, (\text{A/m})$ in
medium 1 (y > 0), determine \mathbf{H}_2 in medium 2 (y < 0).

Solution:

$$\mathbf{J}_{s} = \mathbf{\hat{x}} 8 \text{ A/m},$$
$$\mathbf{H}_{1} = \mathbf{\hat{z}} 11 \text{ A/m}.$$

 \mathbf{H}_1 is tangential to the boundary, and therefore \mathbf{H}_2 is also. With $\hat{\mathbf{n}}_2 = \hat{\mathbf{y}}$, we have

$$\begin{aligned} \hat{\mathbf{n}}_2 \times (\mathbf{H}_1 - \mathbf{H}_2) &= \mathbf{J}_s, \\ \hat{\mathbf{y}} \times (\hat{\mathbf{z}} \, 11 - \mathbf{H}_2) &= \hat{\mathbf{x}} 8, \\ \hat{\mathbf{x}} \, 11 - \hat{\mathbf{y}} \times \mathbf{H}_2 &= \hat{\mathbf{x}} 8, \end{aligned}$$

or

$$\hat{\mathbf{y}} \times \mathbf{H}_2 = \hat{\mathbf{x}} \mathbf{3},$$



Fig04: Adjacent magnetic media with J_s on boundary.

which implies that \mathbf{H}_2 does not have an *x*-component. Also, since $\mu_1 H_{1y} = \mu_2 H_{2y}$ and \mathbf{H}_1 does not have a *y*-component, it follows that \mathbf{H}_2 does not have a *y*-component either. Consequently, we conclude that

 $\mathbf{H}_2 = \hat{\mathbf{z}} \mathbf{3}.$

Problem/3< In Fig.''5, """ the plane defined by x - y = 1 separates medium 1 of permeability μ_1 from medium 2 of permeability μ_2 . If no surface current exists on the boundary and

$$\mathbf{B}_1 = \hat{\mathbf{x}}2 + \hat{\mathbf{y}}3 \quad (\mathrm{T}),$$

find \mathbf{B}_2 and then evaluate your result for $\mu_1 = 5\mu_2$. Hint: Start out by deriving the equation for the unit vector normal to the given plane.

Solution: We need to find $\hat{\mathbf{n}}_2$. To do so, we start by finding any two vectors in the plane x - y = 1, and to do that, we need three non-collinear points in that plane. We choose (0, -1, 0), (1, 0, 0), and (1, 0, 1).

Vector A_1 is from (0, -1, 0) to (1, 0, 0):

$$\mathbf{A}_1 = \hat{\mathbf{x}} \mathbf{1} + \hat{\mathbf{y}} \mathbf{1}.$$

Vector A_2 is from (1,0,0) to (1,0,1):

$$A_2 = \hat{z} 1.$$

Hence, if we take the cross product $A_2 \times A_1$, we end up in a direction normal to the given plane, from medium 2 to medium 1,

$$\hat{\mathbf{n}}_{2} = \frac{\mathbf{A}_{2} \times \mathbf{A}_{1}}{|\mathbf{A}_{2} \times \mathbf{A}_{1}|} = \frac{\hat{\mathbf{z}}_{1} \times (\hat{\mathbf{x}}_{1} + \hat{\mathbf{y}}_{1})}{|\mathbf{A}_{2} \times \mathbf{A}_{1}|} = \frac{\hat{\mathbf{y}}_{1} - \hat{\mathbf{x}}_{1}}{\sqrt{1+1}} = \frac{\hat{\mathbf{y}}}{\sqrt{2}} - \frac{\hat{\mathbf{x}}}{\sqrt{2}}$$



Fig.3: """ Magnetic media separated by the plane x - y = 1

In medium 1, normal component is

$$B_{1n} = \hat{\mathbf{n}}_2 \cdot \mathbf{B}_1 = \left(\frac{\hat{\mathbf{y}}}{\sqrt{2}} - \frac{\hat{\mathbf{x}}}{\sqrt{2}}\right) \cdot (\hat{\mathbf{x}} 2 + \hat{\mathbf{y}} 3) = \frac{3}{\sqrt{2}} - \frac{2}{\sqrt{2}} = \frac{1}{\sqrt{2}},$$
$$\mathbf{B}_{1n} = \hat{\mathbf{n}}_2 B_{1n} = \left(\frac{\hat{\mathbf{y}}}{\sqrt{2}} - \frac{\hat{\mathbf{x}}}{\sqrt{2}}\right) \cdot \frac{1}{\sqrt{2}} = \frac{\hat{\mathbf{y}}}{2} - \frac{\hat{\mathbf{x}}}{2}.$$

Tangential component is

$$\mathbf{B}_{1t} = \mathbf{B}_1 - \mathbf{B}_{1n} = (\hat{\mathbf{x}} \, 2 + \hat{\mathbf{y}} \, 3) - \left(\frac{\hat{\mathbf{y}}}{2} - \frac{\hat{\mathbf{x}}}{2}\right) = \hat{\mathbf{x}} \, 2.5 + \hat{\mathbf{y}} \, 2.5.$$

Boundary conditions:

$$B_{1n} = B_{2n}$$
, or $B_{2n} = \frac{\hat{\mathbf{y}}}{2} - \frac{\hat{\mathbf{x}}}{2}$,
 $H_{1t} = H_{2t}$, or $\frac{B_{2t}}{\mu_2} = \frac{B_{1t}}{\mu_1}$.

Hence,

$$\mathbf{B}_{2t} = \frac{\mu_2}{\mu_1} \mathbf{B}_{1t} = \frac{\mu_2}{\mu_1} (\hat{\mathbf{x}} 2.5 + \hat{\mathbf{y}} 2.5).$$

Finally,

$$\mathbf{B}_2 = \mathbf{B}_{2n} + \mathbf{B}_{2t} = \left(\frac{\hat{\mathbf{y}}}{2} - \frac{\hat{\mathbf{x}}}{2}\right) + \frac{\mu_2}{\mu_1} \left(\hat{\mathbf{x}} 2.5 + \hat{\mathbf{y}} 2.5\right).$$

For $\mu_1 = 5\mu_2$,

$$\mathbf{B}_2 = \hat{\mathbf{y}} \quad (\mathbf{T}).$$

Inductance and Magnetic Energy

Problem/3< Obtain an expression for the self-inductance per unit length for the parallel wire transmission line of Fig.'3(a) in terms of *a*, *d*, and μ , where *a* is the radius of the wires, *d* is the axis-to-axis distance between the wires, and μ is the permeability of the medium in which they reside.

Solution: Let us place the two wires in the x-z plane and orient the current in one of them to be along the +z-direction and the current in the other one to be along the -z-direction, the magnetic field at point P(x,0,z) due to wire 1 is

$$\mathbf{B}_1 = \hat{\mathbf{\phi}} \frac{\mu I}{2\pi r} = \hat{\mathbf{y}} \frac{\mu I}{2\pi x},$$

where the permeability has been generalized from free space to any substance with permeability μ , and it has been recognized that in the *x*-*z* plane, $\hat{\mathbf{\phi}} = \hat{\mathbf{y}}$ and r = x as long as x > 0.

Given that the current in wire 2 is opposite that in wire 1, the magnetic field created by wire 2 at point P(x, 0, z) is in the same direction as that created by wire 1, and it is given by

$$\mathbf{B}_2 = \mathbf{\hat{y}} \frac{\mu I}{2\pi (d-x)}.$$



Fig: Parallel wire transmission line.

Therefore, the total magnetic field in the region between the wires is

$$\mathbf{B} = \mathbf{B}_1 + \mathbf{B}_2 = \hat{\mathbf{y}} \frac{\mu I}{2\pi} \left(\frac{1}{x} + \frac{1}{d-x} \right) = \hat{\mathbf{y}} \frac{\mu I d}{2\pi x (d-x)}.$$

the flux crossing the surface area between the wires over a length l of the wire structure is

$$\Phi = \iint_{S} \mathbf{B} \cdot d\mathbf{s} = \int_{z=z_{0}}^{z_{0}+l} \int_{x=a}^{d-a} \left(\hat{\mathbf{y}} \frac{\mu I d}{2\pi x (d-x)} \right) \cdot \left(\hat{\mathbf{y}} \, dx \, dz \right)$$
$$= \frac{\mu I l d}{2\pi} \left(\frac{1}{d} \ln \left(\frac{x}{d-x} \right) \right) \Big|_{x=a}^{d-a}$$
$$= \frac{\mu I l}{2\pi} \left(\ln \left(\frac{d-a}{a} \right) - \ln \left(\frac{a}{d-a} \right) \right)$$
$$= \frac{\mu I l}{2\pi} \times 2 \ln \left(\frac{d-a}{a} \right) = \frac{\mu I l}{\pi} \ln \left(\frac{d-a}{a} \right).$$

Since the number of 'turns' in this structure is 1, the flux linkage is the same as magnetic flux: $\Lambda = \Phi$. 'y g total inductance over the length *l* ku'i kxgp'as<

$$L = \frac{\Lambda}{I} = \frac{\Phi}{I} = \frac{\mu l}{\pi} \ln\left(\frac{d-a}{a}\right) \quad (\mathrm{H}).$$

Therefore, the inductance per unit length is

$$L' = \frac{L}{l} = \frac{\mu}{\pi} \ln\left(\frac{d-a}{a}\right) \approx \frac{\mu}{\pi} \ln\left(\frac{d}{a}\right)$$
 (H/m),

where the last approximation recognizes that the wires are thin compared to the separation distance (i.e., that $d \gg a$). This has been an implied condition from the beginning of this analysis, where the flux passing through the wires themselves have been ignored.

Problem/4< In terms of the d-c current I, how much magnetic energy is stored in the insulating medium of a 3-m-long, air-filled section of a coaxial transmission line, given that the radius of the inner conductor is 5 cm and the inner radius of the outer conductor is 10 cm?

Solution: the inductance per unit length of an air-filled coaxial cable is given by

$$L' = \frac{\mu_0}{2\pi} \ln\left(\frac{b}{a}\right)$$
 (H/m).

Over a length of 2 m, the inductance is

$$L = 2L' = \frac{3 \times 4\pi \times 10^{-7}}{2\pi} \ln\left(\frac{10}{5}\right) = 416 \times 10^{-9} \quad \text{(H)}.$$

Frqo """" $W_m = LI_I^2 2 = 208I^2$ (nJ), where W_m is in nanojoules when I is in amperes. Alternatively, we can use Eq. (5.106) to compute W_m :

$$W_{\rm m} = \frac{1}{2} \int_{\mathcal{V}} \mu_0 H^2 \, d\mathcal{V}.$$

From Equcvkqp", $H = B/\mu_0 = I/2\pi r$, and

$$W_m = \frac{1}{2} \int_{z=0}^{3m} \int_{\phi=0}^{2\pi} \int_{r=a}^{b} \mu_0 \left(\frac{I}{2\pi r}\right)^2 r \, dr \, d\phi \, dz = 208I^2 \quad \text{(nJ)}.$$