Tutorial#2

Chapter2

Sections-1: Charge and Current Distributions

Problem.1 A cube 2 m on a side is located in the f rst octant in a Cartesian coordinate system, with one of its corners at the origin. Find the total charge contained in the cube if the charge density is given by $\rho_v = xy^2 e^{-2z}$ (mC/m³).

Solution: For the cube shown in Fig. P.1,

$$Q = \int_{\mathcal{V}} \rho_{\nu} d\nu = \int_{x=0}^{2} \int_{y=0}^{2} \int_{z=0}^{2} xy^{2} e^{-2z} dx dy dz$$
$$= \left(\frac{-1}{12}x^{2}y^{3}e^{-2z}\right) \Big|_{x=0}^{2} \Big|_{y=0}^{2} \Big|_{z=0}^{2} = \frac{8}{3}(1 - e^{-4}) = 2.62 \text{ mC}$$



Figure P.1: Cube of Problem.1

Problem.2 Find the total charge contained in a cylindrical volume defined by $r \le 2$ m and $0 \le z \le 3$ m if $\rho_v = 20rz$ (mC/m³).

Solution: For the cylinder shown in Fig. P.2,

$$Q = \int_{z=0}^{3} \int_{\phi=0}^{2\pi} \int_{r=0}^{2} 20rz \, r \, dr \, d\phi \, dz$$
$$= \left(\frac{10}{3}r^{3}\phi z^{2}\right) \Big|_{r=0}^{2} \Big|_{\phi=0}^{2\pi} \Big|_{z=0}^{3} = 480\pi \, (\text{mC}) = 1.5 \, \text{C}.$$



Figure P.2: Cylinder of Problem.2

Problem.3 Find the total charge contained in a cone defined by $R \le 2$ m and $0 \le \theta \le \pi/4$, given that $\rho_v = 10R^2 \cos^2 \theta$ (mC/m³).

Solution: For the cone of Fig. P.3,

$$Q = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi/4} \int_{R=0}^{2} 10R^2 \cos^2 \theta R^2 \sin \theta \, dR \, d\theta \, d\phi$$

= $\left(\frac{-2}{3}R^5 \phi \cos^3 \theta\right) \Big|_{R=0}^{2} \Big|_{\theta=0}^{\pi/4} \Big|_{\phi=0}^{2\pi}$
= $\frac{128\pi}{3} \left(1 - \left(\frac{\sqrt{2}}{2}\right)^3\right) = 86.65$ (mC).



Figure P.3: Cone of Problem.3

Problem.4 If the line charge density is given by $\rho_l = 24y^2$ (mC/m), f nd the total charge distributed on the *y*-axis from y = -5 to y = 5.

Solution:

$$Q = \int_{-5}^{5} \rho_l \, dy = \int_{-5}^{5} 24y^2 \, dy = \left. \frac{24y^3}{3} \right|_{-5}^{5} = 2000 \text{ mC} = 2 \text{ C}.$$

Problem.5 Find the total charge on a circular disk defined by $r \le a$ and z = 0 if:

- (a) $\rho_{\rm s} = \rho_{\rm s0} \cos \phi \, ({\rm C/m^2}),$ (b) $\rho_{\rm s} = \rho_{\rm s0} \sin^2 \phi \, ({\rm C/m^2}),$ (c) $\rho_{\rm s} = \rho_{\rm s0} e^{-r} \, ({\rm C/m^2}),$ (d) $\rho_{\rm s} = \rho_{\rm s0} e^{-r} \sin^2 \phi \, ({\rm C/m^2}),$
- where ρ_{s0} is a constant.

Solution:

(a)

$$Q = \int \rho_{\rm s} \, ds = \int_{r=0}^{a} \int_{\phi=0}^{2\pi} \rho_{\rm s0} \cos\phi \, r \, dr \, d\phi = \rho_{\rm s0} \, \left. \frac{r^2}{2} \right|_{0}^{a} \sin\phi \Big|_{0}^{2\pi} = 0$$

(b)

$$Q = \int_{r=0}^{a} \int_{\phi=0}^{2\pi} \rho_{s0} \sin^{2} \phi \ r \, dr \, d\phi = \rho_{s0} \frac{r^{2}}{2} \Big|_{0}^{a} \int_{0}^{2\pi} \left(\frac{1 - \cos 2\phi}{2}\right) \, d\phi$$
$$= \frac{\rho_{s0} a^{2}}{4} \left(\phi - \frac{\sin 2\phi}{2}\right) \Big|_{0}^{2\pi} = \frac{\pi a^{2}}{2} \rho_{s0}.$$

(c)

$$Q = \int_{r=0}^{a} \int_{\phi=0}^{2\pi} \rho_{s0} e^{-r} r \, dr \, d\phi = 2\pi \rho_{s0} \int_{0}^{a} r e^{-r} \, dr$$
$$= 2\pi \rho_{s0} \left[-r e^{-r} - e^{-r} \right]_{0}^{a}$$
$$= 2\pi \rho_{s0} [1 - e^{-a} (1 + a)]$$

(d)

$$Q = \int_{r=0}^{a} \int_{\phi=0}^{2\pi} \rho_{s0} e^{-r} \sin^{2} \phi \ r \, dr \, d\phi$$

= $\rho_{s0} \int_{r=0}^{a} r e^{-r} \, dr \int_{\phi=0}^{2\pi} \sin^{2} \phi \, d\phi$
= $\rho_{s0} [1 - e^{-a} (1 + a)] \cdot \pi = \pi \rho_{s0} [1 - e^{-a} (1 + a)].$

Problem.6 If $\mathbf{J} = \hat{\mathbf{y}} 4xz$ (A/m²), f nd the current *I* f owing through a square with corners at (0,0,0), (2,0,0), (2,0,2), and (0,0,2).

Solution: the net current f owing through the square shown in Fig. P.6 is

$$I = \int_{S} \mathbf{J} \cdot d\mathbf{s} = \int_{x=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{y}} 4xz) \bigg|_{y=0} \cdot (\hat{\mathbf{y}} \, dx \, dz) = (x^{2}z^{2}) \bigg|_{x=0}^{2} \bigg|_{z=0}^{2} = 16 \text{ A}.$$



Figure P.6: Square surface.

Problem.7 If $\mathbf{J} = \hat{\mathbf{R}}5/R$ (A/m²), f nd *I* through the surface R = 5 m. **Solution:** Using Eq. (4.12), we have

$$I = \int_{S} \mathbf{J} \cdot d\mathbf{s} = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \left(\hat{\mathbf{R}} \frac{5}{R} \right) \cdot \left(\hat{\mathbf{R}} R^{2} \sin \theta \, d\theta \, d\phi \right)$$
$$= -5R\phi \cos \theta \bigg|_{R=5} \bigg|_{\theta=0}^{\pi} \bigg|_{\phi=0}^{2\pi} = 100\pi = 314.2 \quad (A).$$

Problem.8 An electron beam shaped like a circular cylinder of radius r_0 carries a charge density given by

$$\rho_{\rm v} = \left(\frac{-\rho_0}{1+r^2}\right) \quad ({\rm C}/{\rm m}^3), \label{eq:rho_v}$$

where ρ_0 is a positive constant and the beam's axis is coincident with the *z*-axis.

- (a) Determine the total charge contained in length L of the beam.
- (b) If the electrons are moving in the +z-direction with uniform speed u, determine the magnitude and direction of the current crossing the z-plane.

Solution:

(a)

$$Q = \int_{r=0}^{r_0} \int_{z=0}^{L} \rho_{\rm v} \, d\, \nu = \int_{r=0}^{r_0} \int_{z=0}^{L} \left(\frac{-\rho_0}{1+r^2}\right) 2\pi r \, dr \, dz$$
$$= -2\pi\rho_0 L \int_0^{r_0} \frac{r}{1+r^2} \, dr = -\pi\rho_0 L \ln(1+r_0^2).$$

(b)

$$\begin{aligned} \mathbf{J} &= \rho_{\mathbf{v}} \mathbf{u} = -\hat{\mathbf{z}} \frac{u\rho_0}{1+r^2} \quad (A/m^2), \\ I &= \int \mathbf{J} \cdot d\mathbf{s} \\ &= \int_{r=0}^{r_0} \int_{\phi=0}^{2\pi} \left(-\hat{\mathbf{z}} \frac{u\rho_0}{1+r^2} \right) \cdot \hat{\mathbf{z}} r \, dr \, d\phi \\ &= -2\pi u \rho_0 \int_0^{r_0} \frac{r}{1+r^2} \, dr = -\pi u \rho_0 \ln(1+r_0^2) \quad (A). \end{aligned}$$

Current direction is along $-\hat{z}$.

Section-2: Coulomb's Law

Problem.9 A square with sides 2 m each has a charge of $40 \,\mu\text{C}$ at each of its four corners. Determine the electric f eld at a point 5 m above the center of the square.



Figure P.9: Square with charges at the corners.

Solution: The distance |R| between any of the charges and point P is

$$\begin{aligned} |\mathbf{R}| &= \sqrt{1^2 + 1^2 + 5^2} = \sqrt{27}. \\ \mathbf{E} &= \frac{Q}{4\pi\epsilon_0} \left[\frac{\mathbf{R}_1}{|\mathbf{R}|^3} + \frac{\mathbf{R}_2}{|\mathbf{R}|^3} + \frac{\mathbf{R}_3}{|\mathbf{R}|^3} + \frac{\mathbf{R}_4}{|\mathbf{R}|^3} \right] \\ &= \frac{Q}{4\pi\epsilon_0} \left[\frac{-\hat{\mathbf{x}} - \hat{\mathbf{y}} + \hat{\mathbf{z}}5}{(27)^{3/2}} + \frac{\hat{\mathbf{x}} - \hat{\mathbf{y}} + \hat{\mathbf{z}}5}{(27)^{3/2}} + \frac{-\hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}}5}{(27)^{3/2}} + \frac{\hat{\mathbf{x}} + \hat{\mathbf{y}} + \hat{\mathbf{z}}5}{(27)^{3/2}} \right] \\ &= \hat{\mathbf{z}} \frac{5Q}{(27)^{3/2}\pi\epsilon_0} = \hat{\mathbf{z}} \frac{5 \times 40 \ \mu\text{C}}{(27)^{3/2}\pi\epsilon_0} = \frac{1.42}{\pi\epsilon_0} \times 10^{-6} \ (\text{V/m}) = \hat{\mathbf{z}}51.2 \ (\text{kV/m}). \end{aligned}$$

Problem.10 Three point charges, each with q = 3 nC, are located at the corners of a triangle in the *x*-*y* plane, with one corner at the origin, another at (2 cm, 0, 0), and the third at (0, 2 cm, 0). Find the force acting on the charge located at the origin.

Solution: the electric f eld at the origin due to the other two point charges [Fig. P.10]:

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \left[\frac{3 \text{ nC} (-\hat{\mathbf{x}} 0.02)}{(0.02)^3} \right] + \frac{3 \text{ nC} (-\hat{\mathbf{y}} 0.02)}{(0.02)^3} = -67.4(\hat{\mathbf{x}} + \hat{\mathbf{y}}) \text{ (kV/m) at } \mathbf{R} = 0.$$

the force $\mathbf{F} = q\mathbf{E} = -202.2(\hat{\mathbf{x}} + \hat{\mathbf{y}}) (\mu N).$



Figure P.10: Locations of charges in Problem.10.

Problem.11 Charge $q_1 = 6 \ \mu C$ is located at $(1 \ cm, 1 \ cm, 0)$ and charge q_2 is located at $(0, 0, 4 \ cm)$. What should q_2 be so that **E** at $(0, 2 \ cm, 0)$ has no *y*-component?

Solution: For the conf guration of Fig. P.11,



Figure P.11: Locations of charges in Problem.11.

$$\mathbf{E}(\mathbf{R} = \hat{\mathbf{y}}2\text{cm}) = \frac{1}{4\pi\epsilon} \left[\frac{6\mu\text{C}(-\hat{\mathbf{x}} + \hat{\mathbf{y}}) \times 10^{-2}}{(2 \times 10^{-2})^{3/2}} + \frac{q_2(\hat{\mathbf{y}}2 - \hat{\mathbf{z}}4) \times 10^{-2}}{(20 \times 10^{-2})^{3/2}} \right]$$

= $\frac{1}{4\pi\epsilon} [-\hat{\mathbf{x}}21.21 \times 10^{-6} + \hat{\mathbf{y}}(21.21 \times 10^{-6} + 0.224q_2) - \hat{\mathbf{z}}0.447q_2]$ (V/m).

If $E_y = 0$, then $q_2 = -21.21 \times 10^{-6} / 0.224 \approx -94.69 \ (\mu \text{C})$.

Problem.12 A line of charge with uniform density $\rho_l = 8 \ (\mu C/m)$ exists in air along the *z*-axis between z = 0 and z = 5 cm. Find **E** at (0,10 cm,0).

Solution: for the line of charge shown in Fig. P.12 gives

$$\begin{split} \mathbf{E} &= \frac{1}{4\pi\varepsilon_0} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_l \, dl'}{R'^2} \,, \\ R' &= \hat{\mathbf{y}} 0.1 - \hat{\mathbf{z}} z \\ &= \frac{1}{4\pi\varepsilon_0} \int_{z=0}^{0.05} (8 \times 10^{-6}) \frac{(\hat{\mathbf{y}} 0.1 - \hat{\mathbf{z}} z)}{[(0.1)^2 + z^2]^{3/2}} \, dz \\ &= \frac{8 \times 10^{-6}}{4\pi\varepsilon_0} \left[\frac{\hat{\mathbf{y}} 10z + \hat{\mathbf{z}}}{\sqrt{(0.1)^2 + z^2}} \right] \Big|_{z=0}^{0.05} \\ &= 71.86 \times 10^3 \left[\hat{\mathbf{y}} 4.47 - \hat{\mathbf{z}} 1.06 \right] = \hat{\mathbf{y}} \, 321.4 \times 10^3 - \hat{\mathbf{z}} 76.2 \times 10^3 \quad (V/m). \end{split}$$



Figure P.12: Line charge.

Problem.13 Electric charge is distributed along an arc located in the *x*-*y* plane and defined by r = 2 cm and $0 \le \phi \le \pi/4$. If $\rho_l = 5$ (μ C/m), f nd **E** at (0,0,*z*) and then evaluate it at (a) the origin, (b) z = 5 cm, and (c) z = -5 cm.

Solution: For the arc of charge shown in Fig. P.13, $dl = r d\phi = 0.02 d\phi$, and $\mathbf{R}' = -\hat{\mathbf{x}} 0.02 \cos \phi - \hat{\mathbf{y}} 0.02 \sin \phi + \hat{\mathbf{z}} z$.

$$\mathbf{E} = \frac{1}{4\pi\varepsilon_0} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_l \, dl'}{R'^2}$$

= $\frac{1}{4\pi\varepsilon_0} \int_{\phi=0}^{\pi/4} \rho_l \frac{(-\hat{\mathbf{x}}0.02\cos\phi - \hat{\mathbf{y}}0.02\sin\phi + \hat{\mathbf{z}}z)}{((0.02)^2 + z^2)^{3/2}} 0.02 \, d\phi$
= $\frac{898.8}{((0.02)^2 + z^2)^{3/2}} [-\hat{\mathbf{x}}0.014 - \hat{\mathbf{y}}0.006 + \hat{\mathbf{z}}0.78z] \quad (V/m).$

(a) At z = 0, $\mathbf{E} = -\hat{\mathbf{x}} \, 1.6 - \hat{\mathbf{y}} \, 0.66$ (MV/m). (b) At z = 5 cm, $\mathbf{E} = -\hat{\mathbf{x}} \, 81.4 - \hat{\mathbf{y}} \, 33.7 + \hat{\mathbf{z}} \, 226$ (kV/m). (c) At z = -5 cm, $\mathbf{E} = -\hat{\mathbf{x}} \, 81.4 - \hat{\mathbf{y}} \, 33.7 - \hat{\mathbf{z}} \, 226$ (kV/m).

Figure P .13: Line charge along an arc.

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Section-3: Gauss's Law

Problem.14 Three inf nite lines of charge, all parallel to the *z*-axis, are located at the three corners of the kite-shaped arrangement shown in Fig.-14. If the

two right triangles are symmetrical and of equal corresponding sides, show that the electric f eld is zero at the origin.

Figure P.14: Kite-shaped arrangment of line charges for Problem.14.

Solution: The f eld due to an inf nite line of charge is given by the present case, the total **E** at the origin is

$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2 + \mathbf{E}_3.$$

The components of E_1 and E_2 along \hat{x} cancel and their components along $-\hat{y}$ add. Also, E_3 is along \hat{y} because the line charge on the *y*-axis is negative. Hence,

$$\mathbf{E} = -\hat{\mathbf{y}} \frac{2\rho_l \cos\theta}{2\pi\epsilon_0 R_1} + \hat{\mathbf{y}} \frac{2\rho_l}{2\pi\epsilon_0 R_2}$$

But $\cos \theta = R_1/R_2$. Hence,

$$\mathbf{E} = -\hat{\mathbf{y}} \frac{\rho_l}{\pi \varepsilon_0 R_1} \frac{R_1}{R_2} + \hat{\mathbf{y}} \frac{\rho_l}{\pi \varepsilon_0 R_2} = 0.$$

Problem.15 Three inf nite lines of charge, $\rho_{l_1} = 3$ (nC/m), $\rho_{l_2} = -3$ (nC/m), and $\rho_{l_3} = 3$ (nC/m), are all parallel to the *z*-axis. If they pass through the respective points

Figure P.15: Three parallel line charges.

(0,-b), (0,0), and (0,b) in the *x*-*y* plane, f nd the electric f eld at (a,0,0). Evaluate your result for a = 2 cm and b = 1 cm.

Solution:

$$\begin{split} \rho_{l_1} &= 3 \quad (nC/m), \\ \rho_{l_2} &= -3 \quad (nC/m), \\ \rho_{l_3} &= \rho_{l_1}, \\ \mathbf{E} &= \mathbf{E}_1 + \mathbf{E}_2 + \mathbf{E}_3. \end{split}$$

Components of line charges 1 and 3 along y cancel and components along x add. Hence,

$$\mathbf{E} = \hat{\mathbf{x}} \frac{2\rho_{l_1}}{2\pi\epsilon_0 R_1} \cos\theta + \hat{\mathbf{x}} \frac{\rho_{l_2}}{2\pi\epsilon_0 a}.$$

with $\cos \theta = \frac{a}{\sqrt{a^2 + b^2}}$ and $R_1 = \sqrt{a^2 + b^2}$, $\mathbf{E} = \frac{\hat{\mathbf{x}} \, 3}{2\pi\epsilon_0} \left[\frac{2a}{a^2 + b^2} - \frac{1}{a} \right] \times 10^{-9} \quad (V/m).$ For a = 2 cm and b = 1 cm,

$$\mathbf{E} = \hat{\mathbf{x}} \, 1.62 \quad (\mathrm{kV/m}).$$

Problem.16 A horizontal strip lying in the *x*–*y* plane is of width *d* in the *y*-direction and infinitely long in the *x*-direction. If the strip is in air and has a uniform charge distribution ρ_s , use Coulomb's law to obtain an explicit expression for the electric f eld at a point *P* located at a distance *h* above the centerline of the strip. Extend your result to the special case where *d* is infinite.

Figure P.16: Horizontal strip of charge.

Solution: The strip of charge density ρ_s (C/m²) can be treated as a set of adjacent line charges each of charge $\rho_l = \rho_s dy$ and width dy. At point *P*, the f elds of line charge at distance *y* and line charge at distance -y give contributions that cancel each other along $\hat{\mathbf{y}}$ and add along $\hat{\mathbf{z}}$. For each such pair,

$$d\mathbf{E} = \hat{\mathbf{z}} \frac{2\rho_{\rm s} \, dy \cos \theta}{2\pi\varepsilon_0 R}$$

With $R = h/\cos\theta$, we integrate from y = 0 to d/2, which corresponds to $\theta = 0$ to $\theta_0 = \sin^{-1}[(d/2)/(h^2 + (d/2)^2)^{1/2}]$. Thus,

$$\mathbf{E} = \int_0^{d/2} d\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{\pi \varepsilon_0} \int_0^{d/2} \frac{\cos \theta}{R} dy = \hat{\mathbf{z}} \frac{\rho_s}{\pi \varepsilon_0} \int_0^{\theta_0} \frac{\cos^2 \theta}{h} \cdot \frac{h}{\cos^2 \theta} d\theta$$
$$= \hat{\mathbf{z}} \frac{\rho_s}{\pi \varepsilon_0} \theta_0.$$

For an infinitely wide sheet, $\theta_0 = \pi/2$ and $\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0}$,

Problem.17 Given the electric f ux density

$$\mathbf{D} = \hat{\mathbf{x}}2(x+y) + \hat{\mathbf{y}}(3x-2y) \quad (C/m^2),$$

determine

(a) ρ_{v} ,

- (b) the total charge Q enclosed in a cube 2 m on a side, located in the f rst octant with three of its sides coincident with the *x*-, *y*-, and *z*-axes and one of its corners at the origin, and
- (c) the total charge Q in the cube,

Solution:

(a) By applying

$$\rho_{\rm v} = \nabla \cdot \mathbf{D} = \frac{\partial}{\partial x} (2x + 2y) + \frac{\partial}{\partial y} (3x - 2y) = 0.$$

(b) Integrate the charge density over the volume:

$$Q = \int_{\mathcal{V}} \nabla \cdot \mathbf{D} \, dv = \int_{x=0}^{2} \int_{y=0}^{2} \int_{z=0}^{2} 0 \, dx \, dy \, dz = 0.$$

(c) Apply Gauss' law to calculate the total charge

$$Q = \oint \mathbf{D} \cdot d\mathbf{s} = F_{\text{front}} + F_{\text{back}} + F_{\text{right}} + F_{\text{left}} + F_{\text{top}} + F_{\text{bottom}},$$

$$F_{\text{front}} = \int_{y=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{x=2} \cdot (\hat{\mathbf{x}} \, dz \, dy)$$

$$= \int_{y=0}^{2} \int_{z=0}^{2} 2(x+y) \Big|_{x=2} dz \, dy = \left(2z \left(2y + \frac{1}{2}y^{2} \right) \Big|_{z=0}^{2} \right) \Big|_{y=0}^{2} = 24,$$

$$\begin{split} F_{\text{back}} &= \int_{y=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{x=0} \cdot (-\hat{\mathbf{x}} dz dy) \\ &= -\int_{y=0}^{2} \int_{z=0}^{2} 2(x+y) \Big|_{x=0} dz dy = -\left(zy^{2}\Big|_{z=0}^{2}\right)\Big|_{y=0}^{2} = -8, \\ F_{\text{right}} &= \int_{x=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{y=2} \cdot (\hat{\mathbf{y}} dz dx) \\ &= \int_{x=0}^{2} \int_{z=0}^{2} (3x-2y) \Big|_{y=2} dz dx = \left(z\left(\frac{3}{2}x^{2}-4x\right)\Big|_{z=0}^{2}\right)\Big|_{x=0}^{2} = -4, \\ F_{\text{left}} &= \int_{x=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{y=0} \cdot (-\hat{\mathbf{y}} dz dx) \\ &= -\int_{x=0}^{2} \int_{z=0}^{2} (3x-2y) \Big|_{y=0} dz dx = -\left(z\left(\frac{3}{2}x^{2}\right)\Big|_{z=0}^{2}\right)\Big|_{x=0}^{2} = -12, \\ F_{\text{top}} &= \int_{x=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{z=2} \cdot (\hat{z} dy dx) \\ &= \int_{x=0}^{2} \int_{z=0}^{2} 0\Big|_{z=2} dy dx = 0, \\ F_{\text{bottom}} &= \int_{x=0}^{2} \int_{z=0}^{2} (\hat{\mathbf{x}} 2(x+y) + \hat{\mathbf{y}} (3x-2y)) \Big|_{z=0} \cdot (\hat{z} dy dx) \\ &= \int_{x=0}^{2} \int_{z=0}^{2} 0\Big|_{z=2} dy dx = 0. \\ \text{Thus } Q = \oint \mathbf{D} \cdot d\mathbf{s} = 24 - 8 - 4 - 12 + 0 + 0 = 0. \end{split}$$

Problem.18 Repeat Problem.17 for $\mathbf{D} = \hat{\mathbf{x}}xy^3z^3$ (C/m²). Solution:

(a)
$$\rho_{v} = \nabla \cdot \mathbf{D} = \frac{\partial}{\partial x} (xy^{3}z^{3}) = y^{3}z^{3}.$$

(b) Total charge Q is given by:

$$Q = \int_{\mathcal{V}} \nabla \cdot \mathbf{D} \, dv = \int_{z=0}^{2} \int_{y=0}^{2} \int_{x=0}^{2} y^{3}z^{3} \, dx \, dy \, dz = \frac{xy^{4}z^{4}}{16} \Big|_{x=0}^{2} \Big|_{y=0}^{2} \Big|_{z=0}^{2} = 32 \text{ C}.$$

(c) Using Gauss' law we have

$$\oint_{S} \mathbf{D} \cdot d\mathbf{s} = F_{\text{front}} + F_{\text{back}} + F_{\text{right}} + F_{\text{left}} + F_{\text{top}} + F_{\text{bottom}}.$$

Note that $\mathbf{D} = \hat{\mathbf{x}} D_x$, so only F_{front} and F_{back} (integration over $\hat{\mathbf{z}}$ surfaces) will contribute to the integral.

$$F_{\text{front}} = \int_{z=0}^{2} \int_{y=0}^{2} (\hat{\mathbf{x}} x y^{3} z^{3}) \Big|_{x=2} \cdot (\hat{\mathbf{x}} \, dy \, dz)$$

$$= \int_{z=0}^{2} \int_{y=0}^{2} x y^{3} z^{3} \Big|_{x=2} \, dy \, dz = \left(2 \left(\frac{y^{4} z^{4}}{16} \right) \Big|_{y=0}^{2} \right) \Big|_{z=0}^{2} = 32 ,$$

$$F_{\text{back}} = \int_{z=0}^{2} \int_{y=0}^{2} (\hat{\mathbf{x}} x y^{3} z^{3}) \Big|_{x=0} \cdot (-\hat{\mathbf{x}} \, dy \, dz) = -\int_{z=0}^{2} \int_{y=0}^{2} x y^{3} z^{3} \Big|_{x=0} \, dy \, dz = 0.$$

Thus $Q = \oint \mathbf{D} \cdot d\mathbf{s} = 32 + 0 + 0 + 0 + 0 + 0 = 32 \text{ C}.$

Problem.19 Charge Q_1 is uniformly distributed over a thin spherical shell of radius *a*, and charge Q_2 is uniformly distributed over a second spherical shell of radius *b*, with b > a. Apply Gauss's law to f nd **E** in the regions R < a, a < R < b, and R > b.

Solution: Using symmetry considerations, we know $\mathbf{D} = \hat{\mathbf{R}}D_{R}$. From, $d\mathbf{s} = \hat{\mathbf{R}}R^{2}\sin\theta \,d\theta \,d\phi$ for an element of a spherical surface. Using Gauss's law in integral form,

$$\oint_{\mathbf{S}} \mathbf{D} \cdot d\mathbf{s} = Q_{\text{tot}},$$

where Q_{tot} is the total charge enclosed in S. For a spherical surface of radius R,

$$\int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} (\hat{\mathbf{R}} D_{\mathrm{R}}) \cdot (\hat{\mathbf{R}} R^2 \sin \theta \, d\theta \, d\phi) = Q_{\mathrm{tot}},$$
$$D_{\mathrm{R}} R^2 (2\pi) [-\cos \theta]_0^{\pi} = Q_{\mathrm{tot}},$$
$$D_{\mathrm{R}} = \frac{Q_{\mathrm{tot}}}{4\pi R^2}$$

we know a linear, isotropic material has the constitutive relationship $D = \varepsilon E$. Thus, we f nd E from D.

(a) In the region R < a,

$$Q_{\text{tot}} = 0,$$
 $\mathbf{E} = \hat{\mathbf{R}} E_{\text{R}} = \frac{\hat{\mathbf{R}} Q_{\text{tot}}}{4\pi R^2 \varepsilon} = 0$ (V/m).

(b) In the region a < R < b,

$$Q_{\text{tot}} = Q_1, \qquad \mathbf{E} = \hat{\mathbf{R}} E_{\text{R}} = \frac{\hat{\mathbf{R}} Q_1}{4\pi R^2 \varepsilon} \quad (\text{V/m}).$$

(c) In the region R > b,

$$Q_{\text{tot}} = Q_1 + Q_2, \qquad \mathbf{E} = \hat{\mathbf{R}} E_{\text{R}} = \frac{\hat{\mathbf{R}}(Q_1 + Q_2)}{4\pi R^2 \varepsilon} \quad (\text{V/m}).$$

Problem.20 The electric fux density inside a dielectric sphere of radius *a* centered at the origin is given by

$$\mathbf{D} = \hat{\mathbf{R}} \rho_0 R \quad (C/m^2),$$

where ρ_0 is a constant. Find the total charge inside the sphere.

Solution:

$$Q = \oint_{S} \mathbf{D} \cdot d\mathbf{s} = \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \hat{\mathbf{R}} \rho_{0} R \cdot \hat{\mathbf{R}} R^{2} \sin \theta \, d\theta \, d\phi \bigg|_{R=a}$$
$$= 2\pi \rho_{0} a^{3} \int_{0}^{\pi} \sin \theta \, d\theta = -2\pi \rho_{0} a^{3} \cos \theta |_{0}^{\pi} = 4\pi \rho_{0} a^{3} \quad (C).$$

Problem.21 In a certain region of space, the charge density is given in cylindrical coordinates by the function:

$$\rho_{\rm v} = 50 r e^{-r}$$
 (C/m³).

Apply Gauss's law to f nd **D**.

Solution:

Figure P.21: Gaussian surface.

Method 1: Integral Form of Gauss's Law

Since ρ_v varies as a function of *r* only, so will **D**. Hence, we construct a cylinder of radius *r* and length *L*, coincident with the *z*-axis. Symmetry suggests that **D** has the functional form **D** = $\hat{\mathbf{r}}$ *D*. Hence,

$$\oint_{S} \mathbf{D} \cdot d\mathbf{s} = Q,$$

$$\int \hat{\mathbf{r}} D \cdot d\mathbf{s} = D(2\pi rL),$$

$$Q = 2\pi L \int_{0}^{r} 50re^{-r} \cdot r \, dr$$

$$= 100\pi L[-r^{2}e^{-r} + 2(1 - e^{-r}(1 + r))],$$

$$\mathbf{D} = \hat{\mathbf{r}} D = \hat{\mathbf{r}} 50 \left[\frac{2}{r}(1 - e^{-r}(1 + r)) - re^{-r}\right].$$

Method 2: Differential Method

 $\nabla \cdot \mathbf{D} = \rho_{\rm v}, \qquad \mathbf{D} = \hat{\mathbf{r}} D_r,$

with D_r being a function of r.

$$\frac{1}{r}\frac{\partial}{\partial r}(rD_r)=50re^{-r},$$

$$\frac{\partial}{\partial r}(rD_r) = 50r^2e^{-r},$$

$$\int_0^r \frac{\partial}{\partial r}(rD_r) dr = \int_0^r 50r^2e^{-r} dr,$$

$$rD_r = 50[2(1 - e^{-r}(1 + r)) - r^2e^{-r}],$$

$$\mathbf{D} = \hat{\mathbf{r}} rD_r = \hat{\mathbf{r}} 50 \left[\frac{2}{r}(1 - e^{-r}(1 + r)) - re^{-r}\right].$$

Problem.22 An infinitely long cylindrical shell extending between r = 1 m and r = 3 m contains a uniform charge density ρ_{v0} . Apply Gauss's law to f nd **D** in all regions.

Solution: For r < 1 m, $\mathbf{D} = 0$.

For $1 \le r \le 3$ m,

$$\begin{split} \oint_{S} \hat{\mathbf{r}} D_{r} \cdot d\mathbf{s} &= Q, \\ D_{r} \cdot 2\pi rL &= \rho_{v0} \cdot \pi L(r^{2} - 1^{2}), \\ \mathbf{D} &= \hat{\mathbf{r}} D_{r} = \hat{\mathbf{r}} \frac{\rho_{v0} \pi L(r^{2} - 1)}{2\pi rL} = \hat{\mathbf{r}} \frac{\rho_{v0}(r^{2} - 1)}{2r}, \qquad 1 \le r \le 3 \text{ m.} \end{split}$$

For $r \ge 3$ m,

$$D_r \cdot 2\pi r L = \rho_{v0} \pi L (3^2 - 1^2) = 8\rho_{v0} \pi L,$$

$$\mathbf{D} = \hat{\mathbf{r}} D_r = \hat{\mathbf{r}} \frac{4\rho_{v0}}{r}, \qquad r \ge 3 \text{ m}.$$

Figure P.22: Cylindrical shell.

Problem.23 If the charge density increases linearly with distance from the origin such that $\rho_v = 0$ at the origin and $\rho_v = 40 \text{ C/m}^3$ at R = 2 m, f nd the corresponding variation of **D**.

Solution:

$$\rho_{\rm v}(R) = a + bR,$$

$$\rho_{\rm v}(0) = a = 0,$$

$$\rho_{\rm v}(2)=2b=40.$$

Hence, b = 20.

$$\rho_{\rm v}(R) = 20R ~({\rm C/m^3}).$$

Applying Gauss's law to a spherical surface of radius R,

$$\oint_{S} \mathbf{D} \cdot d\mathbf{s} = \int_{V} \rho_{V} d\nu,$$

$$D_{R} \cdot 4\pi R^{2} = \int_{0}^{R} 20R \cdot 4\pi R^{2} dR = 80\pi \frac{R^{4}}{4},$$

$$D_{R} = 5R^{2} \quad (C/m^{2}),$$

$$\mathbf{D} = \hat{\mathbf{R}} D_{R} = \hat{\mathbf{R}} 5R^{2} \quad (C/m^{2}).$$

Section-4: Electric Potential

Problem.24 A square in the x-y plane in free space has a point charge of +Q at corner (a/2, a/2) and the same at corner (a/2, -a/2) and a point charge of -Q at each of the other two corners.

- (a) Find the electric potential at any point *P* along the *x*-axis.
- (b) Evaluate V at x = a/2.

Solution: $R_1 = R_2$ and $R_3 = R_4$.

$$V = \frac{Q}{4\pi\epsilon_0 R_1} + \frac{Q}{4\pi\epsilon_0 R_2} + \frac{-Q}{4\pi\epsilon_0 R_3} + \frac{-Q}{4\pi\epsilon_0 R_4} = \frac{Q}{2\pi\epsilon_0} \left(\frac{1}{R_1} - \frac{1}{R_3}\right)$$

with

$$R_{1} = \sqrt{\left(x - \frac{a}{2}\right)^{2} + \left(\frac{a}{2}\right)^{2}},$$

$$R_{3} = \sqrt{\left(x + \frac{a}{2}\right)^{2} + \left(\frac{a}{2}\right)^{2}}.$$

At x = a/2,

$$R_1 = \frac{a}{2},$$

$$R_3 = \frac{a\sqrt{5}}{2},$$

$$V = \frac{Q}{2\pi\varepsilon_0} \left(\frac{2}{a} - \frac{2}{\sqrt{5}a}\right) = \frac{0.55Q}{\pi\varepsilon_0 a}.$$

Figure P.24: Potential due to four point charges.

Problem.25 The circular disk of radius *a* shown in Fig. (P.25) has uniform charge density ρ_s across its surface.

- (a) Obtain an expression for the electric potential V at a point P(0,0,z) on the z-axis.
- (b) Use your result to f nd E and then evaluate it for z = h. Compare your f nal expression with result which was obtained on the basis of Coulomb's law.

Solution:

(a) Consider a ring of charge at a radial distance r. The charge contained in width dr is

$$dq = \rho_{\rm s}(2\pi r\,dr) = 2\pi\rho_{\rm s}r\,dr.$$

The potential at *P* is

$$dV = \frac{dq}{4\pi\varepsilon_0 R} = \frac{2\pi\rho_{\rm s}r\,dr}{4\pi\varepsilon_0(r^2 + z^2)^{1/2}}.$$

The potential due to the entire disk is

$$V = \int_0^a dV = \frac{\rho_{\rm s}}{2\epsilon_0} \int_0^a \frac{r\,dr}{(r^2 + z^2)^{1/2}} = \frac{\rho_{\rm s}}{2\epsilon_0} \left(r^2 + z^2\right)^{1/2} \bigg|_0^a = \frac{\rho_{\rm s}}{2\epsilon_0} \left[(a^2 + z^2)^{1/2} - z \right].$$

Figure P.25: Circular disk of charge.

(b)

$$\mathbf{E} = -\nabla V = -\hat{\mathbf{x}}\frac{\partial V}{\partial x} - \hat{\mathbf{y}}\frac{\partial V}{\partial y} - \hat{\mathbf{z}}\frac{\partial V}{\partial z} = \hat{\mathbf{z}}\frac{\rho_{s}}{2\varepsilon_{0}}\left[1 - \frac{z}{\sqrt{a^{2} + z^{2}}}\right].$$

 $\sqrt{}$

 \checkmark

The expression for **E** reduces to????????? when z = h.

Problem.26 For the electric dipole shown in Fig. 4-13, d = 1 cm and $|\mathbf{E}| = 4$ (mV/m) at R = 1 m and $\theta = 0^{\circ}$. Find \mathbf{E} at R = 2 m and $\theta = 90^{\circ}$.

Solution: For R = 1 m and $\theta = 0^\circ$, $|\mathbf{E}| = 4$ mV/m, we can solve for *q* using:

$$\mathbf{E} = \frac{qd}{4\pi\varepsilon_0 R^3} (\hat{\mathbf{R}} 2\cos\theta + \hat{\mathbf{\theta}}\sin\theta).$$

Hence,

$$|\mathbf{E}| = \left(\frac{qd}{4\pi\varepsilon_0}\right) 2 = 4 \text{ mV/m} \quad \text{at } \theta = 0^\circ,$$
$$q = \frac{10^{-3} \times 8\pi\varepsilon_0}{d} = \frac{10^{-3} \times 8\pi\varepsilon_0}{10^{-2}} = 0.8\pi\varepsilon_0 \quad (C).$$

Again to f nd **E** at R =

2 m and $\theta = 90^{\circ}$, we have

$$\mathbf{E} = \frac{0.8\pi\varepsilon_0 \times 10^{-2}}{4\pi\varepsilon_0 \times 2^3} (\hat{\mathbf{R}}(0) + \hat{\mathbf{\theta}}) = \hat{\mathbf{\theta}} \frac{1}{4} \quad (\text{mV/m}).$$

Problem.27 For each of the following distributions of the electric potential V, sketch the corresponding distribution of **E** (in all cases, the vertical axis is in volts and the horizontal axis is in meters):

Solution:

Figure P.27: Electric potential distributions

Problem.28 Given the electric f eld

$$\mathbf{E} = \hat{\mathbf{R}} \frac{18}{R^2} \quad (\text{V/m}),$$

f nd the electric potential of point A with respect to point B where A is at +2 m and B at -4 m, both on the z-axis.

Solution:

$$V_{AB} = V_A - V_B = -\int_B^A \mathbf{E} \cdot d\mathbf{l}.$$

Along z-direction, $\hat{\mathbf{R}} = \hat{\mathbf{z}}$ and $\mathbf{E} = \hat{\mathbf{z}}\frac{18}{z^2}$ for $z \ge 0$, and $\hat{\mathbf{R}} = -\hat{\mathbf{z}}$ and $\mathbf{E} = -\hat{\mathbf{z}}\frac{18}{z^2}$ for $z \le 0$. Hence,

$$V_{AB} = -\int_{-4}^{2} \hat{\mathbf{R}} \frac{18}{z^{2}} \cdot \hat{\mathbf{z}} dz = -\left[\int_{-4}^{0} -\hat{\mathbf{z}} \frac{18}{z^{2}} \cdot \hat{\mathbf{z}} dz + \int_{0}^{2} \hat{\mathbf{z}} \frac{18}{z^{2}} \cdot \hat{\mathbf{z}} dz\right] = 4 \text{ V}.$$

Figure P.28: Potential between B and A.

Problem.29 An inf nitely long line of charge with uniform density $\rho_l = 9$ (nC/m) lies in the *x*-*y* plane parallel to the *y*-axis at x = 2 m. Find the potential V_{AB} at point A(3 m, 0, 4 m) in Cartesian coordinates with respect to point B(0, 0, 0)

Solution: According to,

$$V = \frac{\rho_l}{2\pi\varepsilon_0} \ln\left(\frac{r_2}{r_1}\right)$$

where r_1 and r_2 are the distances of A and B. In this case,

$$r_1 = \sqrt{(3-2)^2 + 4^2} = \sqrt{17} \text{ m},$$

 $r_2 = 2 \text{ m}.$

Hence,

$$V_{AB} = \frac{9 \times 10^{-9}}{2\pi \times 8.85 \times 10^{-12}} \ln\left(\frac{2}{\sqrt{17}}\right) = -117.09 \text{ V}.$$

Figure P.29: Line of charge parallel to y-axis.

Problem.30 The *x*-*y* plane contains a uniform sheet of charge with $\rho_{s_1} = 0.2$ (nC/m²) and a second sheet with $\rho_{s_2} = -0.2$ (nC/m²) occupies the plane z = 6 m. Find V_{AB} , V_{BC} , and V_{AC} for A(0,0,6 m), B(0,0,0), and C(0,-2 m,2 m).

Solution: We start by f nding the **E** f eld in the region between the plates. For any point above the x-y plane, **E**₁ due to the charge on x-y plane is,,

$$\mathbf{E}_1 = \hat{\mathbf{z}} \frac{\mathbf{\rho}_{s_1}}{2\mathbf{\epsilon}_0}.$$

In the region below the top plate, **E** would point downwards for positive ρ_{s_2} on the top plate. In this case, $\rho_{s_2} = -\rho_{s_1}$. Hence,

$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2 = \hat{\mathbf{z}} \frac{\rho_{s_1}}{2\epsilon_0} - \hat{\mathbf{z}} \frac{\rho_{s_2}}{2\epsilon_0} = \hat{\mathbf{z}} \frac{2\rho_{s_1}}{2\epsilon_0} = \hat{\mathbf{z}} \frac{\rho_{s_1}}{\epsilon_0}$$

Since **E** is along \hat{z} , only change in position along z can result in change in voltage.

$$V_{AB} = -\int_0^6 \hat{\mathbf{z}} \frac{\rho_{s_1}}{\varepsilon_0} \cdot \hat{\mathbf{z}} dz = -\frac{\rho_{s_1}}{\varepsilon_0} z \Big|_0^6 = -\frac{6\rho_{s_1}}{\varepsilon_0} = -\frac{6 \times 0.2 \times 10^{-9}}{8.85 \times 10^{-12}} = -135.59 \text{ V}.$$

Figure P.30: Two parallel planes of charge.

The voltage at C depends only on the *z*-coordinate of C. Hence, with point A being at the lowest potential and B at the highest potential,

$$V_{BC} = \frac{-2}{6} V_{AB} = -\frac{(-135.59)}{3} = 45.20 \text{ V},$$

$$V_{AC} = V_{AB} + V_{BC} = -135.59 + 45.20 = -90.39 \text{ V}.$$